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Generalized Glaisher-Kinkelin constants, Blagouchine's integrals, and Ramanujan summation of series

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Abstract The main motivation for this article is to establish a connection between a sequence of constants of Glaisher-Kinkelin type (known as the Bendersky-Adamchik constants) which appear quite naturally in number theory, and a family of complex integrals recently introduced by Blagouchine. Furthermore, we also elucidate the close relation between these constants and the Ramanujan summation of certain divergent series, and we give a new convergent series expansion for the logarithm of these constants.

Keywords Generalized Glaisher-Kinkelin constants; generalized gamma functions; infinite series with zeta values; Blagouchine's integrals; Ramanujan summation of series.

1 Introduction

Introduced in 1933 by Bendersky [2], the sequence of mathematical constants called generalized Glaisher-Kinkelin constants or Bendersky-Adamchik constants (see Definition 1) arises quite naturally in number theory [10, 12, 14]. The initial purpose of this article is to link this sequence of real numbers $\{A_k\}$ to the family of complex integrals $\{\mathcal{J}_{k,p}\}$ defined (for nonnegative integers k and p , with p odd) by

$$\mathcal{J}_{k,p} := \int_{-\infty}^{+\infty} \frac{\zeta\left(\frac{p}{2} + ix\right)}{(2k + p + 2ix) \cosh(\pi x)} dx.$$

In an unpublished note, Blagouchine [3] introduced these integrals in the cases $p = 1$ and $p = 3$ only. This connection is deduced from the results of a previous study [8] on one hand (see Lemma 1), and from the residue theorem on the other

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hand (see Proposition 1). In particular, we obtain a nice expression of the integral $\mathcal{J}_k = \mathcal{J}_{k,1}$ in terms of the generalized Glaisher-Kinkelin constants A_j for $0 \leq j \leq k-1$ (see Corollary 1). More precisely, we show that for all $k \geq 1$,

$$\mathcal{J}_k = \int_{-\infty}^{+\infty} \frac{\zeta(\frac{1}{2} + ix)}{(2k+1+2ix) \cosh(\pi x)} dx = \sum_{j=0}^{k-1} (-1)^j \binom{k}{j} \ln A_j - \frac{1}{k} - \frac{1}{(k+1)^2}.$$

Furthermore, we completely elucidate the close relation between the logarithm of the generalized Glaisher-Kinkelin constant A_k and the sum of the divergent series $\sum_{n \geq 1} n^k \ln n$ in the sense of Ramanujan's summation method, following the exposition in [4]. More precisely, using the notations of [4], we show that for all $k \geq 0$,

$$\ln A_k = \sum_{n \geq 1}^{\mathcal{R}} n^k \ln n + \frac{H_k B_{k+1}}{k+1} + \frac{1}{(k+1)^2},$$

where, in this expression, H_k and B_k denote respectively the k th harmonic number and the k th Bernoulli number.

Finally, we present, in the last section, a new convergent series expansion for the constant $\ln A_k$ involving the Bernoulli numbers of the second kind (see Corollary 2).

2 Constants of Glaisher-Kinkelin type and Ramanujan summation

Definition 1. Let $\{B_{2r}\}_{r \geq 0}$ be the sequence of (even) Bernoulli numbers. The sequence of numbers $\{A_k\}_{k \geq 0}$ can be defined as follows [14, Eq. (1.1)–(1.6)]: for all $k \geq 0$,

$$\ln A_k := \lim_{n \rightarrow \infty} \left\{ \sum_{\nu=1}^n \nu^k \ln \nu - P_k(n) \ln n + Q_k(n) \right\},$$

where P_k and Q_k are polynomials of degree $k+1$ whose expressions are given by

$$P_0(x) = x + \frac{1}{2}, Q_0(x) = x,$$

and for $k \geq 1$,

$$P_k(x) = \frac{x^{k+1}}{k+1} + \frac{x^k}{2} + \sum_{r=1}^{\lfloor \frac{k+1}{2} \rfloor} \frac{B_{2r}}{(2r)!} \left(\prod_{j=1}^{2r-1} (k-j+1) \right) x^{k+1-2r},$$

$$Q_k(x) = \frac{x^{k+1}}{(k+1)^2} - \sum_{r=1}^{\lfloor \frac{k+1}{2} \rfloor + \frac{(-1)^{k-1}}{2}} \frac{B_{2r}}{(2r)!} \left\{ \prod_{j=1}^{2r-1} (k-j+1) \sum_{j=1}^{2r-1} \frac{1}{k-j+1} \right\} x^{k+1-2r}.$$

The numbers A_k (for $k = 0, 1, 2, \dots$) are called the *generalized Glaisher-Kinkelin constants* or the *Bendersky-Adamchik constants*. Adamchik [1, Prop. 4] has given a nice expression of these constants A_k in terms of the derivatives of the Riemann zeta function. More precisely, this expression (called Adamchik's formula in the remainder of this article) is the following:

$$\ln A_k = \frac{H_k B_{k+1}}{k+1} - \zeta'(-k), \quad (1)$$

where H_k and B_k denote respectively the k th harmonic number (with the usual convention $H_0 = 0$) and the k th Bernoulli number.

Example 1. a) The constant A_0 is the Stirling constant

$$A_0 = \sqrt{2\pi} = \lim_{n \rightarrow \infty} \left\{ \frac{n!}{n^{n+\frac{1}{2}} e^{-n}} \right\} = \exp(-\zeta'(0)).$$

b) The constant A_1 is the classical Glaisher-Kinkelin constant [1, 10, 14]. We have

$$A_1 = \lim_{n \rightarrow \infty} \left\{ \frac{\prod_{\nu=1}^n \nu^\nu}{n^{\frac{n^2}{2} + \frac{n}{2} + \frac{1}{12}} e^{-\frac{n^2}{4}}} \right\} = \exp \left\{ \frac{1}{12} - \zeta'(-1) \right\}.$$

c) For $k = 2$, we have

$$A_2 = \lim_{n \rightarrow \infty} \left\{ \frac{\prod_{\nu=1}^n \nu^{\nu^2}}{n^{\frac{n^3}{3} + \frac{n^2}{2} + \frac{n}{6}} e^{-\frac{n^3}{9} + \frac{n}{12}}} \right\} = \exp(-\zeta'(-2)).$$

Remark 1. Bendersky [2] introduced for the first time the sequence of numbers A_k without any consideration of their relation with the ζ -function. From the point of view of the summation of divergent series, the constants $\ln A_k$ can be interpreted as follows: for any non-negative integer k , if $\sum_{n \geq 1}^{\mathcal{R}} n^k \ln n$ denotes the \mathcal{R} -sum of the divergent series $\sum_{n \geq 1} n^k \ln n$ (i.e. the sum of the series in the sense of Ramanujan's summation method, following the notations of [4]), then $\ln A_k$ and $\sum_{n \geq 1}^{\mathcal{R}} n^k \ln n$ are linked by the simple relation:

$$\ln A_k = \sum_{n \geq 1}^{\mathcal{R}} n^k \ln n + \frac{H_k B_{k+1}}{k+1} + \frac{1}{(k+1)^2} \quad (k \geq 0). \quad (2)$$

This relation results both from [4, p. 68] and from Adamchik's formula (1). By [2, Eq. (V_k), p. 280], we also deduce the equivalent relation

$$\sum_{n \geq 1}^{\mathcal{R}} n^k \ln n = \ln A_k - \frac{H_k B_{k+1}}{k+1} - \frac{1}{(k+1)^2} = \int_0^1 \ln \Gamma_k(x+1) dx, \quad (2 \text{ b})$$

where Γ_k is the Bendersky generalized gamma function [2, pp. 279–280]. This function verifies

$$\Gamma_k(n+1) = \prod_{\nu=1}^n \nu^{\nu^k} \text{ for any integer } n \geq 1,$$

and

$$\Gamma_k(n+1) \sim A_k n^{P_k(n)} e^{-Q_k(n)} \text{ as } n \rightarrow \infty.$$

In particular, $\Gamma_0 = \Gamma$, and Γ_1 is the Kinkelin-Bendersky hyperfactorial function [13, Def. 3]. Unaware of Bendersky's work and following an idea of Milnor, Kurokawa and Ochiai [11, Thm. 2] provided an expression of the function Γ_k in terms of the derivative of the Hurwitz zeta function $\zeta(s, x)$ at $s = -k$. Precisely, they established the relation¹

$$\ln \Gamma_k(x) = \zeta'(-k, x) - \zeta'(-k) \quad (x > 0, k \geq 0),$$

which is a generalization of the classical formula for Γ [7, Def. 9.6.13]

$$\ln \Gamma(x) = \zeta'(0, x) - \zeta'(0) \quad (x > 0).$$

Remark 2. The following identities:

$$\zeta'(-2k) = (-1)^k \frac{(2k)!}{2(2\pi)^{2k}} \zeta(2k+1) \quad (k \geq 1),$$

and

$$\zeta'(1-2k) = (-1)^{k+1} \frac{(2k)!}{k(2\pi)^{2k}} \zeta'(2k) + \frac{B_{2k}}{2k} (H_{2k-1} - \gamma - \ln 2\pi) \quad (k \geq 1),$$

where γ denotes the Euler-Mascheroni constant, are easily derived (by differentiation) from the functional equation of the zeta function. Thanks to Adamchik's formula (1) and the celebrated Euler formula

$$\zeta(2k) = (-1)^{k+1} \frac{(2\pi)^{2k}}{2(2k)!} B_{2k} \quad (k \geq 1),$$

we can deduce the following two ones:

$$\ln A_{2k-1} = \frac{B_{2k}}{2k} \left(\gamma + \ln 2\pi - \frac{\zeta'(2k)}{\zeta(2k)} \right) \quad (k \geq 1), \quad (3)$$

1. According to Kellner [10, Rem. 27], in the special case where $x = n + 1$ is an integer, this expression of $\ln \Gamma_k$ is due to Alexeiewsky.

and

$$\ln A_{2k} = \frac{B_{2k}}{4} \left(\frac{\zeta(2k+1)}{\zeta(2k)} \right) \quad (k \geq 1). \quad (4)$$

These identities will be useful in the remainder (see Example 3 below).

Example 2. The simplest cases of these identities are the following:

$$\ln A_1 = \frac{1}{12} \left(\gamma + \ln 2\pi - \frac{\zeta'(2)}{\zeta(2)} \right) \quad \text{and} \quad \ln A_2 = \frac{1}{24} \left(\frac{\zeta(3)}{\zeta(2)} \right).$$

In particular, for small values of k , formula (2 b) translates into

$$\begin{aligned} \sum_{n \geq 1}^{\mathcal{R}} \ln n &= \frac{1}{2} \ln 2\pi - 1, \\ \sum_{n \geq 1}^{\mathcal{R}} n \ln n &= \frac{1}{12} \left(\gamma + \ln 2\pi - \frac{\zeta'(2)}{\zeta(2)} \right) - \frac{1}{3}, \\ \sum_{n \geq 1}^{\mathcal{R}} n^2 \ln n &= \frac{1}{24} \left(\frac{\zeta(3)}{\zeta(2)} \right) - \frac{1}{9}. \end{aligned}$$

3 Generalized Glaisher-Kinkelin constants and Blagouchine's integrals

Definition 2. For non-negative integers k and p , with p odd, the integral $\mathcal{J}_{k,p}$ and the series $\mathcal{S}_{k,p}$ are defined respectively by

$$\mathcal{J}_{k,p} := \int_{-\infty}^{+\infty} \frac{\zeta\left(\frac{p}{2} + ix\right)}{(2k + p + 2ix) \cosh(\pi x)} dx,$$

and

$$\mathcal{S}_{k,p} := \sum_{n=N_p}^{\infty} (-1)^n \frac{\zeta(n)}{n+k} \quad \text{with } N_p = \max\left(2, \frac{p+1}{2}\right).$$

To lighten our notations, in the remainder of the text, we set

$$\mathcal{J}_k := \int_{-\infty}^{+\infty} \frac{\zeta\left(\frac{1}{2} + ix\right)}{(2k+1+2ix) \cosh(\pi x)} dx = \mathcal{J}_{k,1},$$

and

$$\mathcal{S}_k := \sum_{n=2}^{\infty} (-1)^n \frac{\zeta(n)}{n+k} = \mathcal{S}_{k,1} = \mathcal{S}_{k,3}.$$

Lemma 1. For all $k \geq 1$, we have

$$\mathcal{S}_k = \frac{\gamma}{k+1} + \frac{1}{k} - \sum_{j=0}^{k-1} (-1)^j \binom{k}{j} \ln A_j. \quad (5)$$

Proof. It results from [8, Prop. 1] that

$$\mathcal{S}_k = \frac{\gamma}{k+1} + \sum_{j=0}^{k-1} (-1)^j \binom{k}{j} \zeta'(-j) + \frac{1}{k} + \sum_{j=0}^{k-1} \binom{k}{j} \frac{B_{j+1} H_j}{j+1} \quad (k \geq 1).$$

Using Adamchik's formula (1), this later expression is equivalent to (5). \square

The following proposition includes and extends the results of [3].

Proposition 1. For all $k \geq 0$, we have the following relations:

$$\mathcal{J}_k = \frac{\gamma}{k+1} - \frac{1}{(k+1)^2} - \mathcal{S}_k, \quad (6)$$

and

$$\mathcal{J}_{k,p} = (-1)^{\frac{p+1}{2}} \mathcal{S}_{k,p} \quad (p = 3, 5, 7, \dots). \quad (7)$$

In particular, $\mathcal{J}_{k,3} = \mathcal{S}_k$.

Proof. For $k \geq 0$, let us consider the function

$$f_k(z) = \frac{\zeta(z)}{(k+z) \sin(\pi z)}.$$

This function f_k has poles at integers $n \in \mathbb{Z}$. For $n \geq 2$, the residue of f_k at $z = n$ is

$$\text{Res}(f_k; n) = \frac{(-1)^n \zeta(n)}{(n+k)\pi}.$$

For $n = 1$, f_k has a double pole and

$$\text{Res}(f_k; 1) = -\frac{1}{\pi} \left(\frac{\gamma}{k+1} - \frac{1}{(k+1)^2} \right).$$

If p, q are positive odd integers with $p < q$, then, by the residue theorem, we have

$$\int_{\text{Re}(z)=p/2} f_k(z) dz - \int_{\text{Re}(z)=q/2} f_k(z) dz = -2i\pi \sum_{\frac{q}{2} > n > \frac{p}{2}} \text{Res}(f_k; n). \quad (*)$$

Moreover, there is a positive constant C such that

$$\left| \int_{\text{Re}(z)=q/2} f_k(z) dz \right| \leq C \int_{-\infty}^{+\infty} \frac{1}{\left((k + \frac{q}{2})^2 + t^2 \right)^{\frac{1}{2}} (e^{\pi t} - e^{-\pi t})} dt,$$

and thus, by the dominated convergence theorem, we have

$$\left| \int_{\operatorname{Re}(z)=q/2} f_k(z) dz \right| \rightarrow 0 \quad \text{as } q \rightarrow +\infty.$$

Therefore, taking the limit in (*), we obtain

$$\int_{\operatorname{Re}(z)=p/2} f_k(z) dz = -2i\pi \sum_{n>\frac{p}{2}} \operatorname{Res}(f_k; n). \quad (**)$$

This last identity allows us to deduce formulas (6) and (7). \square

From Lemma 1 and Proposition 1 above, we derive the following corollary.

Corollary 1. For all $k \geq 1$, we have

$$\mathcal{J}_k = \sum_{j=0}^{k-1} (-1)^j \binom{k}{j} \ln A_j - \frac{1}{k} - \frac{1}{(k+1)^2}. \quad (8)$$

Proof. Formula (8) is a direct consequence of (5) and (6). \square

Example 3. For small positive values of k , we have

$$\begin{aligned} \mathcal{J}_1 &= \frac{1}{2} \ln 2\pi - \frac{5}{4}, \\ \mathcal{J}_2 &= \frac{1}{3} \ln 2\pi - \frac{1}{6} \gamma - \frac{11}{18} + \frac{1}{6} \frac{\zeta'(2)}{\zeta(2)}, \\ \mathcal{J}_3 &= \frac{1}{4} \ln 2\pi - \frac{1}{4} \gamma - \frac{19}{48} + \frac{1}{4} \frac{\zeta'(2)}{\zeta(2)} + \frac{1}{8} \frac{\zeta(3)}{\zeta(2)}. \end{aligned}$$

Moreover, using formulas (3) and (4), we obtain the following general formula:

$$\begin{aligned} \mathcal{J}_k &= \frac{1}{k+1} \ln 2\pi - \frac{(k-1)}{2(k+1)} \gamma - \frac{k^2+3k+1}{k(k+1)^2} \\ &\quad + \sum_{j=1}^{\lfloor \frac{k}{2} \rfloor} \binom{k}{2j-1} \frac{B_{2j} \zeta'(2j)}{2j \zeta(2j)} + \frac{1}{4} \sum_{j=1}^{\lfloor \frac{k-1}{2} \rfloor} \binom{k}{2j} \frac{B_{2j} \zeta(2j+1)}{\zeta(2j)} \quad (k \geq 3). \quad (9) \end{aligned}$$

Remark 3. Since $\mathcal{J}_0 = -1$ (by (6)), we can write (8) in a slightly different manner:

$$\sum_{j=0}^{k-1} (-1)^j \binom{k}{j} \ln A_j = \mathcal{J}_k - \frac{k^2+3k+1}{k(k+1)^2} \mathcal{J}_0 \quad (k \geq 1).$$

Remark 4. By means of a Fourier transform method, Candelpergher [5] has recently established a beautiful relation which is a natural generalization of (6) involving a complex parameter s . More precisely, for any non-negative integer k and any complex number s such that $\operatorname{Re}(s) > \frac{1}{2}$, we have the following relation [5, Eq. (7)]:

$$2^{s-1} \mathcal{J}_k(s) = \frac{\gamma}{(k+1)^s} - \frac{s}{(k+1)^{s+1}} - \mathcal{S}_k(s). \quad (10)$$

with

$$\mathcal{J}_k(s) := \int_{-\infty}^{+\infty} \frac{\zeta(\frac{1}{2} + ix)}{(2k+1 + 2ix)^s \cosh(\pi x)} dx,$$

and

$$\mathcal{S}_k(s) := \sum_{n=2}^{\infty} (-1)^n \frac{\zeta(n)}{(n+k)^s}.$$

In particular, applying (10) to the special case $k = 0$, allows us deduce the following interesting identity:

$$\sum_{n=2}^{\infty} (-1)^n \frac{\zeta(n)}{n^s} = \gamma - s - \frac{1}{2} \int_{-\infty}^{+\infty} \frac{\zeta(\frac{1}{2} + ix)}{(\frac{1}{2} + ix)^s \cosh(\pi x)} dx \quad (\operatorname{Re}(s) > \frac{1}{2}). \quad (11)$$

4 New series expansions for the logarithm of the generalized Glaisher-Kinkelin constants

In his seminal work of 1933, Bendersky [2, p. 295–299] had presented two convergent series expansions of the logarithm $L_k = \ln A_k$ of the generalized Glaisher-Kinkelin constants. In this last section, we give a new one, of a different kind, involving the Bernoulli numbers of the second kind.

We first define a sequence of positive rational numbers $\{\lambda_n\}_{n \geq 1}$ (called non-alternating Cauchy numbers [6]) by

$$\lambda_n := \left| \sum_{k=1}^n \frac{s(n, k)}{k+1} \right| \quad (n \geq 1),$$

where $s(n, k)$ denotes the (signed) Stirling numbers of the first kind. Alternatively, these numbers can also be defined recursively by means of the relation

$$\sum_{k=1}^{n-1} \frac{\lambda_k}{k! (n-k)} = \frac{1}{n} \quad (n \geq 2).$$

The first ones are the following:

$$\lambda_1 = \frac{1}{2}, \lambda_2 = \frac{1}{6}, \lambda_3 = \frac{1}{4}, \lambda_4 = \frac{19}{30}, \lambda_5 = \frac{9}{4}, \lambda_6 = \frac{863}{84}, \text{ etc.}$$

The numbers λ_n are closely related to the Bernoulli numbers of the second kind b_n [9] through the relation

$$\lambda_n = (-1)^{n-1} n! b_n = n! |b_n| \quad (n \geq 1).$$

We now give another interesting application of Adamchik's formula (1).

Proposition 2. Let $S(k, r)$ be the Stirling numbers of the second kind

$$S(k, r) = \frac{1}{r!} \sum_{j=0}^r (-1)^{r-j} \binom{r}{j} j^k \quad (0 \leq r \leq k),$$

and σ_r the shifted Mascheroni series

$$\sigma_r := \sum_{n=r+1}^{\infty} \frac{\lambda_n}{n!(n-r)} \quad (r \geq 0).$$

Then, for all integers $k \geq 1$, we have

$$\ln A_k = (-1)^k \sum_{r=1}^k (-1)^r r! S(k, r) \sigma_{r+1} + \frac{B_{k+1}}{k+1} (H_{k+1} + \gamma). \quad (12)$$

Proof. Thanks to Adamchik's formula (1), formula (12) can be easily deduced from the decomposition of $\zeta'(-k)$ given by [9, Prop. 3]. \square

Remark 5. This formula (12) also extends to the case $k = 0$ through the identity:

$$\ln A_0 = \frac{1}{2} \ln(2\pi) = \sigma_1 + \frac{1}{2} \gamma + \frac{1}{2} = \sum_{n=2}^{\infty} \frac{\lambda_n}{n!(n-1)} + \frac{1}{2} (1 + \gamma).$$

Corollary 2. For all integers $k \geq 1$, we have

$$\ln A_{2k} = \sum_{n=2k+2}^{\infty} \frac{\lambda_n}{n!} \left\{ \sum_{r=1}^{2k} \frac{(-1)^r r! S(2k, r)}{n-1-r} \right\} + C_{2k}, \quad (13)$$

and

$$\ln A_{2k-1} = \sum_{n=2k+1}^{\infty} \frac{\lambda_n}{n!} \left\{ \sum_{r=1}^{2k-1} \frac{(-1)^{r-1} r! S(2k-1, r)}{n-1-r} \right\} + \frac{B_{2k}}{2k} (H_{2k} + \gamma) + C_{2k-1}, \quad (14)$$

where the constants C_k are given by $C_1 = 0$, and

$$C_k = (-1)^k \sum_{r=1}^{k-1} (-1)^r r! S(k, r) \sum_{j=r+2}^{k+1} \frac{\lambda_j}{j!(j-1-r)} \quad (k \geq 2).$$

Example 4.

$$\begin{aligned}\ln A_1 &= \sum_{n=3}^{\infty} \frac{\lambda_n}{n!(n-2)} + \frac{1}{12}\gamma + \frac{1}{8} \\ \ln A_2 &= \sum_{n=4}^{\infty} \frac{\lambda_n(n-1)}{n!(n-2)(n-3)} - \frac{1}{24} \\ \ln A_3 &= \sum_{n=5}^{\infty} \frac{\lambda_n n(n-1)}{n!(n-2)(n-3)(n-4)} - \frac{1}{120}\gamma - \frac{29}{240} \\ \ln A_4 &= \sum_{n=6}^{\infty} \frac{\lambda_n(n-1)^2(n+4)}{n!(n-2)(n-3)(n-4)(n-5)} - \frac{113}{480} \\ \ln A_5 &= \sum_{n=7}^{\infty} \frac{\lambda_n n(n-1)(n^2+13n-18)}{n!(n-2)(n-3)(n-4)(n-5)(n-6)} + \frac{1}{252}\gamma - \frac{55087}{80640}\end{aligned}$$

Remark 6. Thanks to the simple relation linking $\ln A_k$ to the \mathcal{R} -sum $\sum_{n \geq 1}^{\mathcal{R}} n^k \ln n$ given by (2), we can also deduce from the previous corollary, the corresponding formulas:

$$\sum_{n \geq 1}^{\mathcal{R}} n^{2k} \ln n = \sum_{n=2k+2}^{\infty} \frac{\lambda_n}{n!} \left\{ \sum_{r=1}^{2k} \frac{(-1)^r r! S(2k, r)}{n-1-r} \right\} - \frac{1}{(2k+1)^2} + C_{2k}, \quad (15)$$

$$\sum_{n \geq 1}^{\mathcal{R}} n^{2k-1} \ln n = \sum_{n=2k+1}^{\infty} \frac{\lambda_n}{n!} \left\{ \sum_{r=1}^{2k-1} \frac{(-1)^{r-1} r! S(2k-1, r)}{n-1-r} \right\} + \frac{B_{2k}}{2k} \gamma + D_{2k}, \quad (16)$$

with

$$D_{2k} = \frac{B_{2k} - 1}{(2k)^2} + C_{2k-1} \quad (k \geq 1).$$

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