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# Generalized Glaisher-Kinkelin constants and Blagouchine's integrals 

Marc-Antoine Coppo*<br>Université Côte d'Azur, CNRS, LJAD (UMR 7351), Nice, France


#### Abstract

The main purpose of this article is to establish a close connection between two sequences of complex integrals introduced by Blagouchine and some important mathematical constants, namely the Euler-Mascheroni constant, the Cohen-Boyadzhiev constant, and the generalized Glaisher-Kinkelin constants (also known as the Bendersky constants) which occur quite naturally in analysis and number theory. At the end of this study, we also use a formula of Candelpergher to deduce an interesting expression of the alternating series $\sum_{n \geq 2}(-1)^{n} \frac{\zeta(n)}{n^{s}}$ for complex values of $s$.


Keywords Generalized Glaisher-Kinkelin constants, infinite series with zeta values, complex integration.

## 1 Introduction

The main purpose of this article is to highlight the link between the sequence of complex integrals $\left\{\mathcal{I}_{k}\right\}_{k \geq 0}$ and $\left\{\mathcal{J}_{k}\right\}_{k \geq 0}$ respectively defined by

$$
\mathcal{I}_{k}=\int_{-\infty}^{+\infty} \frac{\zeta\left(\frac{3}{2}+i x\right)}{(2 k+1+2 i x) \cosh (\pi x)} d x
$$

and

$$
\mathcal{J}_{k}=\int_{-\infty}^{+\infty} \frac{\zeta\left(\frac{1}{2}+i x\right)}{(2 k+1+2 i x) \cosh (\pi x)} d x
$$

on one side, and some important mathematical constants, namely the EulerMascheroni constant, the Cohen-Boyadzhiev constant, and the generalized GlaisherKinkelin constants (also known as the Bendersky constants) which occur quite

[^0]naturally in analysis and number theory [7, 10, 11]. To establish these close connections, we make use of a relation (found by Blagouchine [2, Theorem 1]) between the integral $\mathcal{J}_{k}$ and the alternating series
$$
\nu_{k}:=\sum_{n=2}^{\infty}(-1)^{n} \frac{\zeta(n)}{n+k}
$$
studied in [8]. Recently, this deep relation has been thoroughly generalized by Candelpergher [6, Eq. (7)], that allows us to give an interesting expression of the alternating series
$$
\sum_{n=2}^{\infty}(-1)^{n} \frac{\zeta(n)}{n^{s}}
$$
for any complex number $s$ with $\operatorname{Re}(s)>\frac{1}{2}$ (see formula (11)).

## 2 Generalized Glaisher-Kinkelin constants

Definition 1 ( $[1,10,11])$. For any integer $k \geq 0$, the constant $A_{k}$ is usually defined by

$$
\ln \left(A_{k}\right)=\lim _{N \rightarrow \infty}\left\{\sum_{n=1}^{N} n^{k} \ln n-P_{k}(N)\right\},
$$

where $P_{k}(N)$ is given by $P_{0}(N)=\left(N+\frac{1}{2}\right) \ln N-N$, and

$$
\begin{aligned}
& P_{k}(N)=\left(\frac{N^{k+1}}{k+1}+\frac{N^{k}}{2}+k!\sum_{j=1}^{k} \frac{N^{k-j} B_{j+1}}{(j+1)!(k-j)!}\right) \ln N \\
& \quad-\frac{N^{k+1}}{(k+1)^{2}}+k!\sum_{j=1}^{k} \frac{N^{k-j} B_{j+1}}{(j+1)!(k-j)!}\left\{\left(1-\delta_{k, j}\right) \sum_{i=1}^{j} \frac{1}{k-i+1}\right\} \quad(k \geq 1),
\end{aligned}
$$

where $B_{j}$ is the $j$-th Bernoulli number and $\delta_{k, j}$ is the Kronecker delta function. The numbers $A_{k}$ (for $k=0,1,2, \ldots$ ) are the generalized Glaisher-Kinkelin constants (sometimes called the Bendersky constants). In particular, it follows from this definition that

$$
\ln \left(A_{0}\right)=\lim _{N \rightarrow \infty}\left\{\sum_{n=1}^{N} \ln n-\left(N+\frac{1}{2}\right) \ln N+N\right\}
$$

and

$$
\ln \left(A_{1}\right)=\lim _{N \rightarrow \infty}\left\{\sum_{n=1}^{N} n \ln n-\left(\frac{N^{2}}{2}+\frac{N}{2}+\frac{1}{12}\right) \ln N+\frac{N^{2}}{4}\right\} .
$$

Remark 1. Adamchik [1, Proposition 4] has found a nice expression of theses constants in terms of the derivatives of the Riemann zeta function. More precisely, we have

$$
\begin{equation*}
A_{k}=\exp \left\{\frac{H_{k} B_{k+1}}{k+1}-\zeta^{\prime}(-k)\right\} \quad(k \geq 0) \tag{1}
\end{equation*}
$$

where $H_{k}=\sum_{j=1}^{k} \frac{1}{j}$ is the $k$-th harmonic number with the usual convention $H_{0}=0$. The following relations are easily deduced by differentiation of Riemann's functional equation for the zeta function:

$$
\zeta^{\prime}(-2 k)=(-1)^{k} \frac{(2 k)!}{2(2 \pi)^{2 k}} \zeta(2 k+1) \quad(k \geq 1)
$$

and

$$
\zeta^{\prime}(1-2 k)=(-1)^{k+1} \frac{(2 k)!}{k(2 \pi)^{2 k}} \zeta^{\prime}(2 k)+\frac{B_{2 k}}{2 k}\left(H_{2 k-1}-\gamma-\ln 2 \pi\right) \quad(k \geq 1) .
$$

This enable to deduce from Adamchik's formula (1) the expressions

$$
\begin{equation*}
A_{2 k-1}=\exp \left\{(-1)^{k} \frac{(2 k)!}{k(2 \pi)^{2 k}} \zeta^{\prime}(2 k)+\frac{B_{2 k}}{2 k}(\gamma+\ln 2 \pi)\right\} \quad(k \geq 1) \tag{2}
\end{equation*}
$$

and

$$
\begin{equation*}
A_{2 k}=\exp \left\{(-1)^{k+1} \frac{(2 k)!}{2(2 \pi)^{2 k}} \zeta(2 k+1)\right\} \quad(k \geq 1) . \tag{3}
\end{equation*}
$$

Example 1. The constant

$$
A_{0}=\exp \left(-\zeta^{\prime}(0)\right)=\sqrt{2 \pi}
$$

is the Stirling constant,

$$
A_{1}=\exp \left(\frac{1}{12}-\zeta^{\prime}(-1)\right)=\exp \left(-\frac{\zeta^{\prime}(2)}{2 \pi^{2}}+\frac{\ln (2 \pi)+\gamma}{12}\right)
$$

is the classical Glaisher-Kinkelin constant, and for $k=2$, we have

$$
A_{2}=\exp \left(-\zeta^{\prime}(-2)\right)=\exp \left(\frac{\zeta(3)}{4 \pi^{2}}\right)
$$

Remark 2. Bendersky [3] introduced for the first time the sequence of numbers $L_{k}:=\ln \left(A_{k}\right)$ without any consideration of their relation with the $\zeta$-function. From the point of view of the summation of divergent series, the constants $L_{k}$ can be interpreted as follows: if $\sum_{n>1}^{\mathcal{R}} n^{k} \ln n$ denotes the $\mathcal{R}$-sum of the divergent series $\sum_{n \geq 1} n^{k} \ln n$ (i.e. the sum of the series in the sense of Ramanujan's summation method [5]), then, for any integer $k \geq 0$, we have

$$
\begin{equation*}
\sum_{n \geq 1}^{\mathcal{R}} n^{k} \ln n=L_{k}-\frac{H_{k} B_{k+1}}{k+1}-\frac{1}{(k+1)^{2}}=\int_{0}^{1} \ln \Gamma_{k}(x+1) d x \tag{4}
\end{equation*}
$$

where $\Gamma_{k}$ is the Bendersky generalized gamma function [3]. This function verifies in particular

$$
\Gamma_{k}(n+1)=1^{1^{k}} 2^{2^{k}} \cdots n^{n^{k}} \quad \text { for any integer } n \geq 1
$$

Apparently unaware of Bendersky's work, Kurokawa and Ochiai [9, Theorem 2] have given a deep expression of the function $\Gamma_{k}$ in terms of the derivative of the Hurwitz zeta function $\zeta(s, x)$ at $s=-k$. Precisely, they showed that

$$
\Gamma_{k}(x)=\exp \left\{\zeta^{\prime}(-k, x)-\zeta^{\prime}(-k)\right\} \quad \text { for } x>0 .
$$

Example 2. For the first values of $k$, formula (4) translates into the identities

$$
\begin{aligned}
& \sum_{n \geq 1}^{\mathcal{R}} \ln n=\ln (\sqrt{2 \pi})-1=\int_{0}^{1} \ln \Gamma(x+1) d x \\
& \sum_{n \geq 1}^{\mathcal{R}} n \ln n=\ln \left(A_{1}\right)-\frac{1}{3}=\int_{0}^{1} \ln K(x+1) d x
\end{aligned}
$$

where $K=\Gamma_{1}$ is the classical Kinkelin hyperfactorial function.

## 3 Blagouchine's first integral

We give a direct proof of [2, Theorem 2] using Cauchy's residue theorem.
Proposition 1. For any integer $k \geq 0$, let $\mathcal{I}_{k}$ be the integral defined by

$$
\mathcal{I}_{k}:=\int_{-\infty}^{+\infty} \frac{\zeta\left(\frac{3}{2}+i x\right)}{(2 k+1+2 i x) \cosh (\pi x)} d x
$$

Then

$$
\begin{equation*}
\mathcal{I}_{k}=\mu_{k} \tag{5}
\end{equation*}
$$

where $\mu_{k}$ is the conditionally convergent series defined by

$$
\mu_{k}:=\sum_{n=1}^{\infty}(-1)^{n-1} \frac{\zeta(n+1)}{n+k} .
$$

Proof. For $k \geq 0$, let us consider the function

$$
f_{k}(z)=\frac{\zeta\left(\frac{3}{2}+i z\right)}{\left(\frac{1}{2}+k+i z\right) \cosh (\pi z)}
$$

We have $\cosh (\pi z)=0$ if and only if $z=i / 2+i n$ with $n \in \mathbb{Z}$. For $n \geq 1$, the residue of $f_{k}$ at $z=i / 2-i n$ is

$$
\frac{\zeta(1+n)}{(n+k) \pi \sinh \left(i \pi\left(\frac{1}{2}-n\right)\right)}=\frac{\zeta(1+n)}{(n+k) i \pi \sin \left(\pi\left(\frac{1}{2}-n\right)\right)}=\frac{(-1)^{n} \zeta(1+n)}{(n+k) i \pi} .
$$

We integrate on a closed contour composed of the interval $D_{R}=[-R, R]$ and the lower semicircle $C_{R}$ of radius $R$ with center at 0 . By the residue theorem, we can then write the following relation:

$$
\frac{1}{2 i \pi} \int_{C_{R}} f_{k}(z) d z+\frac{1}{2 i \pi} \int_{D_{R}} f_{k}(z) d z=-\sum_{n=1}^{N_{R}} \operatorname{Res}\left(f_{k} ; \frac{i}{2}-i n\right)
$$

which, from the foregoing, translates into the identity

$$
\begin{equation*}
\int_{C_{R}} f_{k}(z) d z+\int_{D_{R}} f_{k}(z) d z=2 \sum_{n=1}^{N_{R}}(-1)^{n+1} \frac{\zeta(1+n)}{(n+k)} . \tag{6}
\end{equation*}
$$

For $z \in C_{R}$, the parameterization $i z=R e^{i t}$ with $-\pi / 2<t<\pi / 2$, enables us to write

$$
\begin{array}{r}
\left|\int_{C_{R}} f_{k}(z) d z\right|=\left|\int_{-\pi / 2}^{+\pi / 2} \frac{\zeta\left(\frac{3}{2}+R e^{i t}\right)}{\left(\frac{1}{2}+k+R e^{i t}\right) \cosh \left(i \pi R e^{i t}\right)} R e^{i t} d t\right| \\
\leq \int_{-\pi / 2}^{+\pi / 2}\left|\frac{\zeta\left(\frac{3}{2}+R e^{i t}\right)}{\left(\frac{1}{2}+k+R e^{i t}\right) \cosh \left(i \pi R e^{i t}\right)}\right| R d t .
\end{array}
$$

Since $\frac{3}{2}+R e^{i t}$ is in the half-plane $\operatorname{Re}(z)>3 / 2$, its absolute value is bounded by $\zeta\left(\frac{3}{2}\right)$, i.e.

$$
\left|\zeta\left(\frac{3}{2}+R e^{i t}\right)\right| \leq \zeta\left(\frac{3}{2}\right) .
$$

Hence, when $R$ increases towards infinity, we have the following limits:

$$
\begin{gathered}
\lim _{R \rightarrow \infty} \int_{C_{R}} f_{k}(z) d z=0 \\
\lim _{R \rightarrow \infty} \int_{D_{R}} f_{k}(z) d z=\int_{-\infty}^{+\infty} \frac{\zeta\left(\frac{3}{2}+i x\right)}{\left(\frac{1}{2}+k+i x\right) \cosh (\pi x)} d x=2 \mathcal{I}_{k},
\end{gathered}
$$

and

$$
\lim _{R \rightarrow \infty} \sum_{n=1}^{N_{R}}(-1)^{n+1} \frac{\zeta(1+n)}{n+k}=\sum_{n=1}^{\infty}(-1)^{n+1} \frac{\zeta(n+1)}{n+k}=\mu_{k} .
$$

This allows us to deduce formula (5) by passing to the limit in (6).

Remark 3. The constant

$$
\mu_{0}=\sum_{n=2}^{\infty}(-1)^{n} \frac{\zeta(n)}{n-1}=\sum_{n=1}^{\infty} \frac{1}{n} \ln \left(1+\frac{1}{n}\right)=1.257746 \ldots
$$

has been thoroughly studied by Boyadzhiev [4] (see also [7, p. 142]). This constant is noted $M$ in [4], $K$ in [7], and also appears in [8] as $\nu_{-1}$. By a well-known series representation of Euler's constant $\gamma$, we also have

$$
\mu_{1}=\sum_{n=2}^{\infty}(-1)^{n} \frac{\zeta(n)}{n}=\sum_{n=1}^{\infty}\left(\frac{1}{n}-\ln \left(1+\frac{1}{n}\right)\right)=\gamma=0.577215 \ldots .
$$

Example 3. For $k=0$, formula (5) translates into

$$
\mathcal{I}_{0}=\int_{-\infty}^{+\infty} \frac{\zeta\left(\frac{3}{2}+i x\right)}{(1+2 i x) \cosh (\pi x)} d x=\int_{0}^{1} \frac{\psi(x+1)+\gamma}{x} d x
$$

where $\psi$ is the digamma function, and for $k=1$

$$
\mathcal{I}_{1}=\int_{-\infty}^{+\infty} \frac{\zeta\left(\frac{3}{2}+i x\right)}{(3+2 i x) \cosh (\pi x)} d x=\gamma=-\psi(1)
$$

## 4 Blagouchine's second integral

Proposition 2. For any integer $k \geq 0$, let $\mathcal{J}_{k}$ be the integral defined by

$$
\mathcal{J}_{k}:=\int_{-\infty}^{+\infty} \frac{\zeta\left(\frac{1}{2}+i x\right)}{(2 k+1+2 i x) \cosh (\pi x)} d x
$$

Then

$$
\begin{equation*}
\mathcal{J}_{k}=\frac{\gamma}{k+1}-\frac{1}{(k+1)^{2}}-\mathcal{I}_{k+1} \quad(k \geq 0) . \tag{7}
\end{equation*}
$$

Since $\mathcal{I}_{1}=\gamma$, it follows from this relation that $\mathcal{J}_{0}=-1$; furthermore we have the identity

$$
\begin{equation*}
\mathcal{J}_{k}=\sum_{j=0}^{k-1}(-1)^{j}\binom{k}{j} \ln \left(A_{j}\right)-\frac{k^{2}+3 k+1}{k(k+1)^{2}} \quad(k \geq 1) . \tag{8}
\end{equation*}
$$

Proof. Formula (7) is a reformulation of [2, Theorem 1]. Now assume that $k \geq 1$ and let $\nu_{k}:=\mu_{k+1}$. By [8, Proposition 1] and (5), we have

$$
\mathcal{I}_{k+1}=\nu_{k}=\frac{\gamma}{k+1}+\sum_{j=0}^{k-1}(-1)^{j}\binom{k}{j} \zeta^{\prime}(-j)+\frac{1}{k}+\sum_{j=0}^{k-1}\binom{k}{j} \frac{B_{j+1} H_{j}}{j+1} .
$$

Now, using Adamchik's formula (1), this expression may be rewritten as follows:

$$
\mathcal{I}_{k+1}=\frac{\gamma}{k+1}+\frac{1}{k}-\sum_{j=0}^{k-1}(-1)^{j}\binom{k}{j} \ln \left(A_{j}\right) \quad(k \geq 1) .
$$

Hence, formula (8) results from (7).
Example 4. For $k=1$, formula (8) simply reduces to

$$
\mathcal{J}_{1}=\int_{-\infty}^{+\infty} \frac{\zeta\left(\frac{1}{2}+i x\right)}{(3+2 i x) \cosh (\pi x)} d x=\frac{1}{2} \ln (2 \pi)-\frac{5}{4}
$$

and for $k=2$,

$$
\mathcal{J}_{2}=\int_{-\infty}^{+\infty} \frac{\zeta\left(\frac{1}{2}+i x\right)}{(5+2 i x) \cosh (\pi x)} d x=\frac{1}{3} \ln (2 \pi)-\frac{1}{6} \gamma+\frac{\zeta^{\prime}(2)}{\pi^{2}}-\frac{11}{19} .
$$

Remark 4. The following general expression can be easily deduced from formulas (2) and (3):

$$
\begin{align*}
\sum_{j=0}^{k-1}(-1)^{j}\binom{k}{j} \ln \left(A_{j}\right)= & \frac{\ln (2 \pi)}{k+1}-\frac{k-1}{k+1} \frac{\gamma}{2} \\
& \quad-\sum_{j=1}^{\left[\frac{k}{2}\right]}(-1)^{j}\binom{k}{2 j-1} \frac{(2 j)!}{j(2 \pi)^{2 j}} \zeta^{\prime}(2 j) \\
& \quad-\sum_{j=1}^{\left[\frac{k-1}{2}\right]}(-1)^{j}\binom{k}{2 j} \frac{(2 j)!}{2(2 \pi)^{2 j}} \zeta(2 j+1) \quad(k \geq 3) . \tag{9}
\end{align*}
$$

## 5 Further generalization

Using Fourier transform method, Candelpergher [6, Eq. (7)] recently proved the following beautiful relation which generalizes (7):

Theorem 1. for $k \geq 0$ and $\operatorname{Re}(s)>\frac{1}{2}$, we have

$$
\begin{equation*}
2^{s-1} \mathcal{J}_{k}(s)=\frac{\gamma}{(k+1)^{s}}-\frac{s}{(k+1)^{s+1}}-\nu_{k}(s), \tag{10}
\end{equation*}
$$

with

$$
\mathcal{J}_{k}(s):=\int_{-\infty}^{+\infty} \frac{\zeta\left(\frac{1}{2}+i x\right)}{(2 k+1+2 i x)^{s} \cosh (\pi x)} d x
$$

and

$$
\nu_{k}(s):=\sum_{n=2}^{\infty}(-1)^{n} \frac{\zeta(n)}{(n+k)^{s}} .
$$

Applying (10) with $k=0$ allows us to deduce the following identity:

## Corollary 1.

$$
\begin{equation*}
\sum_{n=2}^{\infty}(-1)^{n} \frac{\zeta(n)}{n^{s}}=\gamma-s-\frac{1}{2} \int_{-\infty}^{+\infty} \frac{\zeta\left(\frac{1}{2}+i x\right)}{\left(\frac{1}{2}+i x\right)^{s} \cosh (\pi x)} d x \quad\left(\operatorname{Re}(s)>\frac{1}{2}\right) \tag{11}
\end{equation*}
$$

Example 5.

$$
\sum_{n=2}^{\infty}(-1)^{n} \frac{\zeta(n)}{n^{2}}=\gamma-2-\frac{1}{2} \int_{-\infty}^{+\infty} \frac{\zeta\left(\frac{1}{2}+i x\right)}{\left(\frac{1}{2}+i x\right)^{2} \cosh (\pi x)} d x
$$

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[^0]:    *Corresponding author. Email address: coppo@unice.fr

